

Chapter 38

2. Let

$$\frac{1}{2}m_e v^2 = E_{\text{photon}} = \frac{hc}{\lambda}$$

and solve for v :

$$\begin{aligned} v &= \sqrt{\frac{2hc}{\lambda m_e}} = \sqrt{\frac{2hc}{\lambda m_e c^2}} c^2 = c \sqrt{\frac{2hc}{\lambda (m_e c^2)}} \\ &= (2.998 \times 10^8 \text{ m/s}) \sqrt{\frac{2(1240 \text{ eV} \cdot \text{nm})}{(590 \text{ nm})(511 \times 10^3 \text{ eV})}} = 8.6 \times 10^5 \text{ m/s}. \end{aligned}$$

Since $v \ll c$, the nonrelativistic formula $K = \frac{1}{2}mv^2$ may be used. The $m_e c^2$ value of Table 37-3 and $hc = 1240 \text{ eV} \cdot \text{nm}$ are used in our calculation.

4. We denote the diameter of the laser beam as d . The cross-sectional area of the beam is $A = \pi d^2/4$. From the formula obtained in Problem 38-3, the rate is given by

$$\begin{aligned} \frac{R}{A} &= \frac{\lambda P}{hc(\pi d^2/4)} = \frac{4(633 \text{ nm})(5.0 \times 10^{-3} \text{ W})}{\pi(6.63 \times 10^{-34} \text{ J} \cdot \text{s})(2.998 \times 10^8 \text{ m/s})(3.5 \times 10^{-3} \text{ m})^2} \\ &= 1.7 \times 10^{21} \text{ photons/m}^2 \cdot \text{s}. \end{aligned}$$

12. Following Sample Problem — “Emission and absorption of light as photons,” we have

$$P = \frac{Rhc}{\lambda} = \frac{(100/\text{s})(6.63 \times 10^{-34} \text{ J} \cdot \text{s})(2.998 \times 10^8 \text{ m/s})}{550 \times 10^{-9} \text{ m}} = 3.6 \times 10^{-17} \text{ W}.$$

20. Using the value $hc = 1240 \text{ eV} \cdot \text{nm}$, the number of photons emitted from the laser per unit time is

$$R = \frac{P}{E_{\text{ph}}} = \frac{2.00 \times 10^{-3} \text{ W}}{(1240 \text{ eV} \cdot \text{nm} / 600 \text{ nm})(1.60 \times 10^{-19} \text{ J / eV})} = 6.05 \times 10^{15} / \text{s},$$

of which $(1.0 \times 10^{-16})(6.05 \times 10^{15}/\text{s}) = 0.605/\text{s}$ actually cause photoelectric emissions. Thus the current is

$$i = (0.605/\text{s})(1.60 \times 10^{-19} \text{ C}) = 9.68 \times 10^{-20} \text{ A}.$$

26. To find the longest possible wavelength λ_{max} (corresponding to the lowest possible energy) of a photon that can produce a photoelectric effect in platinum, we set $K_{\text{max}} = 0$ in Eq. 38-5 and use $hf = hc/\lambda$. Thus $hc/\lambda_{\text{max}} = \Phi$. We solve for λ_{max} :

$$\lambda_{\text{max}} = \frac{hc}{\Phi} = \frac{1240 \text{ eV} \cdot \text{nm}}{5.32 \text{ eV}} = 233 \text{ nm}.$$

47. If K is given in electron volts, then

$$\begin{aligned} \lambda &= \frac{h}{p} = \frac{h}{\sqrt{2mK}} = \frac{6.626 \times 10^{-34} \text{ J} \cdot \text{s}}{\sqrt{2(9.109 \times 10^{-31} \text{ kg})(1.602 \times 10^{-19} \text{ J/eV})K}} = \frac{1.226 \times 10^{-9} \text{ m} \cdot \text{eV}^{1/2}}{\sqrt{K}} \\ &= \frac{1.226 \text{ nm} \cdot \text{eV}^{1/2}}{\sqrt{K}}, \end{aligned}$$

where K is the kinetic energy. Thus,

$$K = \left(\frac{1.226 \text{ nm} \cdot \text{eV}^{1/2}}{\lambda} \right)^2 = \left(\frac{1.226 \text{ nm} \cdot \text{eV}^{1/2}}{590 \text{ nm}} \right)^2 = 4.32 \times 10^{-6} \text{ eV}.$$

50. (a) The momentum of the electron is

$$p = \frac{h}{\lambda} = \frac{6.63 \times 10^{-34} \text{ J} \cdot \text{s}}{0.20 \times 10^{-9} \text{ m}} = 3.3 \times 10^{-24} \text{ kg} \cdot \text{m/s}.$$

(b) The momentum of the photon is the same as that of the electron:
 $p = 3.3 \times 10^{-24} \text{ kg} \cdot \text{m/s}.$

(c) The kinetic energy of the electron is

$$K_e = \frac{p^2}{2m_e} = \frac{(3.3 \times 10^{-24} \text{ kg} \cdot \text{m/s})^2}{2(9.11 \times 10^{-31} \text{ kg})} = 6.0 \times 10^{-18} \text{ J} = 38 \text{ eV}.$$

(d) The kinetic energy of the photon is

$$K_{\text{ph}} = pc = (3.3 \times 10^{-24} \text{ kg} \cdot \text{m/s})(2.998 \times 10^8 \text{ m/s}) = 9.9 \times 10^{-16} \text{ J} = 6.2 \text{ keV}.$$

63. If the momentum is measured at the same time as the position, then

$$\Delta p \approx \frac{\hbar}{\Delta x} = \frac{6.63 \times 10^{-34} \text{ J} \cdot \text{s}}{2\pi(50 \text{ pm})} = 2.1 \times 10^{-24} \text{ kg} \cdot \text{m/s} .$$

64. (a) Using the value $hc = 1240 \text{ nm} \cdot \text{eV}$, we have

$$E = \frac{hc}{\lambda} = \frac{1240 \text{ nm} \cdot \text{eV}}{10.0 \times 10^{-3} \text{ nm}} = 124 \text{ keV} .$$

(b) The kinetic energy gained by the electron is equal to the energy decrease of the photon:

$$\begin{aligned} \Delta E &= \Delta \left(\frac{hc}{\lambda} \right) = hc \left(\frac{1}{\lambda} - \frac{1}{\lambda + \Delta \lambda} \right) = \left(\frac{hc}{\lambda} \right) \left(\frac{\Delta \lambda}{\lambda + \Delta \lambda} \right) = \frac{E}{1 + \lambda / \Delta \lambda} \\ &= \frac{E}{1 + \frac{\lambda}{\lambda_c (1 - \cos \phi)}} = \frac{124 \text{ keV}}{1 + \frac{10.0 \text{ pm}}{(2.43 \text{ pm})(1 - \cos 180^\circ)}} \\ &= 40.5 \text{ keV} . \end{aligned}$$

(c) It is impossible to “view” an atomic electron with such a high-energy photon, because with the energy imparted to the electron the photon would have knocked the electron out of its orbit.

79. The de Broglie wavelength for the bullet is

$$\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{6.63 \times 10^{-34} \text{ J} \cdot \text{s}}{(40 \times 10^{-3} \text{ kg})(1000 \text{ m/s})} = 1.7 \times 10^{-35} \text{ m} .$$

82. The difference between the electron-photon scattering process in this problem and the one studied in the text (the Compton shift, see Eq. 38-11) is that the electron is in motion relative with speed v to the laboratory frame. To utilize the result in Eq. 38-11, shift to a new reference frame in which the electron is at rest before the scattering. Denote the quantities measured in this new frame with a prime ('), and apply Eq. 38-11 to yield

$$\Delta \lambda' = \lambda' - \lambda'_0 = \frac{h}{m_e c} (1 - \cos \pi) = \frac{2h}{m_e c} ,$$

where we note that $\phi = \pi$ (since the photon is scattered back in the direction of incidence). Now, from the Doppler shift formula (Eq. 38-25) the frequency f'_0 of the photon prior to the scattering in the new reference frame satisfies

$$f'_0 = \frac{c}{\lambda'_0} = f_0 \sqrt{\frac{1+\beta}{1-\beta}},$$

where $\beta = v/c$. Also, as we switch back from the new reference frame to the original one after the scattering

$$f = f' \sqrt{\frac{1-\beta}{1+\beta}} = \frac{c}{\lambda'} \sqrt{\frac{1-\beta}{1+\beta}}.$$

We solve the two Doppler-shift equations above for λ' and λ'_0 and substitute the results into the Compton shift formula for $\Delta\lambda'$:

$$\Delta\lambda' = \frac{1}{f} \sqrt{\frac{1-\beta}{1+\beta}} - \frac{1}{f_0} \sqrt{\frac{1-\beta}{1+\beta}} = \frac{2h}{m_e c^2}.$$

Some simple algebra then leads to

$$E = hf = hf_0 \left(1 + \frac{2h}{m_e c^2} \sqrt{\frac{1+\beta}{1-\beta}} \right)^{-1}.$$